

ISM

Simon Casassus Astronomía, Universidad de Chile http:://www.das.uchile.cl/~simon

- I Introduction: Observations of the ISM
- II Microscopic Processes.
- III Astrophysics of Gaseous Nebulae.

- IV Interstellar Dust.
- V Dynamics of the ISM.
- VI Selectec topics: protoplanetary disks, planetary nebulae, SNRs.

Part II

Microscopic Processes

1 Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density lonic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis Ion-molecule chemistry Observational validation Chemical fractionation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Thermal Equilibrium of the ISM



In steady state, dn/dt = dT/dt = 0, so the temperature *T* is determined by $\Lambda(T) = \Gamma(T)$.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Outline

1 Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

Ionic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Cooling of the ISM

Radiative cooling (ref. Dyson & Williams, Chap. 3, Spitzer, Chap. 4 & 6, Osterbrock, Chap. 3).

$$A + B \rightarrow A + B^{\star}$$

 $B^{\star} \rightarrow B + h\nu$ Requisites for a mechanism of radiative cooling

- frequent encounters, i.e. A & B abundant
- $E(B^{\star}) E(B) \stackrel{<}{\sim} kT.$
- large rate of B excitations: high 'collision strength'
- time scale for radiative decay of B^{\star} less than $\tau_{\rm col}$
- large escape probability of the output $h\nu$.



ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Detailed balance: abundance of excited species

In steady state, the density n_j of species B^* in level *j* is determined by an equation of detailed balance, which models systems with a finite number of microscopic states:

$$\sum_{k} \left(\left\{ \begin{array}{c} n_{\mathrm{e}} \\ n_{\mathrm{H}_{2}} \\ n_{\mathrm{H}_{1}} \end{array} \right\} n_{j} \gamma_{jk} + n_{j} B_{jk} J_{\nu} \right) + \sum_{k < j} (A_{jk} n_{j})$$
$$= \sum_{k} \left(n_{k} \left\{ \begin{array}{c} n_{\mathrm{e}} \\ n_{\mathrm{H}_{2}} \\ n_{\mathrm{H}_{1}} \end{array} \right\} \gamma_{kj} + B_{kj} J_{\nu} n_{k} \right) + \sum_{k > j} (n_{k} A_{kj}),$$

which we may also write using a simplified notation:

$$\sum_{i \neq j} n_j C_{ji} + n_j B_{ji} U_{\nu_{ji}} + \sum_{i < j} n_j A_{ji} = \sum_{i \neq j} n_i C_{ij} + n_i B_{ji} U_{\nu_{ij}} + \sum_{i > j} n_i A_{ij}.$$





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Detailed balance: rate of collisonal excitation

For any collision partner (i.e. either $n_{\rm H_2}$), the specific rate of $n_{\rm H_1}$ collisional excitations is

$$\gamma_{jk} = \langle u\sigma_{jk}(u)
angle = rac{4}{\sqrt{\pi}} \left(rac{m_r}{2kT}
ight)^{3/2} \int_0^\infty du \, u^3 \sigma_{jk}(u) \exp\left(-rac{m_r}{2kT}u^2
ight) du \, u^3 \sigma_{jk}(u)$$

the specific rate (i.e. per pair of particles, in units of m³ s⁻¹) of collisional de-excitations γ_{kj} must be related to γ_{jk} : in thermodynamic equilibrium any microscopic process must be balanced by its inverse (*principle of detailed balance*, otherwise the distribution functions would depend on *t* and there would be no equilibrium).



ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Milne relations

In **LTE** the rate of collisional excitations $j \rightarrow k$ with relative velocity $\vec{u} \in [\vec{u}, \vec{u} + d\vec{u}]$ must equal the rate of de-excitations $k \rightarrow j$ in the corresponding velocity range \vec{v} ,

$$n_{j}^{\text{LTE}} \begin{array}{c} n_{e} & n_{e} \\ n_{H_{2}} & \sigma_{jk}(u)f(u)ud\vec{u} = n_{k}^{\text{LTE}} & n_{H_{2}} & \sigma_{kj}(v)f(v)vd\vec{v}, \\ n_{H_{1}} & n_{H_{1}} \end{array}$$

with
$$f(u) = \left(\frac{1}{\pi} \frac{m}{2kT}\right)^{3/2} \exp\left(-\frac{m}{2kT}u^2\right)$$
,
and $\frac{1}{2}m_r v^2 = \frac{1}{2}m_r u^2 - E_{jk}^{-1}$,

$$\stackrel{ ext{tarea}}{\Longrightarrow} \quad g_j u^2 \sigma_{jk}(u) = g_k v^2 \sigma_{kj}(v), ext{ or }$$
 $C_{kj} = C_{jk} \frac{g_j}{g_k} e^{E_{jk}/kT}.$

ISM



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

¹note threshold excitation energy

Critical density

The *critical density* is defined as the density at which the rate of collisonal de-excitations is equal to the rate of radiative de-excitations:

$$\sum_{k < j} C_{kj} = \sum_{k < j} A_{jk}$$

A good cooling mechanism is therefore one that involves abundant species, an effective collision strength $\gtrsim 1$, a critical density $n_{\rm crit} \gg n_{\rm e}$, and optically thin transport of the output $h\nu_{jk}$.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Outline

1 Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density lonic & atomic cooling of the ISM

Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Ionic & atomic cooling of the ISM

For ionic collisions, the cross-section of collisional excitation highlights the Coulomb cross-section.

$$\sigma_{jk}(E) = rac{\pi a_{\circ}^2 \Omega_{jk}(E)}{E/E_{\circ}},$$

where E_{\circ} is 1 Rydberg, a_{\circ} is the Bohr radius, and $\Omega_{jk}(E)$ (of order 1), is the *collision strength* as a function of kinetic energy *E*. Finally the rate of excitations $j \to k$ (with units of s⁻¹) is (tarea),

$$C_{jk} = \left\{ \begin{array}{c} n_{\rm e} \\ n_{\rm H_2} \\ n_{\rm H_1} \end{array} \right\} \sqrt{\frac{2\pi}{kT}} \left(\frac{h}{2\pi} \right)^2 \frac{1}{m_r^{3/2}} \frac{\Gamma_{jk}}{g_j} e^{-E_{jk}/kT},$$

where Γ_{jk} , the *effective collision strength*, is the Maxwellian average of $\Omega(E)$.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Fine structure cooling in the IR

In LS coupling energy levels are ordered according to the Hund rules (Shu I Chap. 27):

- Higher S \rightarrow lower energy
- Higher L \rightarrow lower energy
- Higher J \rightarrow higher energy if less than half-filled, lower energy if more than half-filled

Spectroscopic notation in spin-orbit coupling (ver Shu I,27):

$${}^{2S+1}L_J$$
 , with degeneracy (2 J+1) Examples:

- [Ar VI] 4.52 μ m, ${}^{2}P_{\frac{1}{2}} \leftarrow {}^{2}P_{\frac{3}{2}}$, Z=18
- [Al vI] 3.65 μ m, ${}^{3}P_{2} \leftarrow {}^{3}P_{1}$, Z=13



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

1.12

Ar VI ${}^{2}P_{\frac{3}{2}}$ population





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

Ar VI ${}^{2}P_{\frac{3}{2}}$ population





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

Cooling of H II regions

For gas in a steady state, where every ionization event is balanced by a recombination event (case of a nebula photoionised by OB stars), the cooling mechanisms are:

- Recombination radiation (postponed to chapter on photoionised nebulae)
- 2 Continuum emission (free-free, bound-free): $A + B \rightarrow A + B + h\nu$ (postponed to chapter on photoionised nebulae)
- 3 Collisionally excited line emission:

$$\Lambda_{CE} = \sum_{X} \sum_{i} n(X, i) \sum_{j < i} A_{i,j} (E_i - E_j)$$

e.g.

- $[O II]\lambda\lambda 3726,3729$ ${}^{4}S_{3/2} \leftarrow {}^{2}D_{3/2,5/2}.$ $[O III]\lambda\lambda 4959,5007$ ${}^{3}P_{2,1} \leftarrow {}^{1}D_{2}.$



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Cooling of HII regions: collisionally excited line radiation





ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

Ionic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

1.16

Cooling of H $\scriptstyle \parallel$ regions: collisionally excited line radiation $z\sim z_{\odot}$





ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

Ionic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Cooling of H $\scriptstyle\rm II$ regions: collisionally excited line radiation $z\sim 0.1 z_{\odot}$



Thermal Equilibrium of the ISM

ISM

Cooling of the ISM: detailed balance, critical density

Ionic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Cooling of H $\scriptstyle\rm II$ regions: collisionally excited line radiation $z\sim 3z_\odot$





ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

Ionic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Cooling of the neutral ISM

Occurs through excitation of low-lying electronic states, fine structure levels dominate.

Examples: Table 3.1 in Dyson & Williams complemented by the list of lines in http://www.pa.uky.edu/~peter/atomic/

Transition	collision partner	$\Delta E/k$
[C I] λ 157.7 μ m ${}^{2}P_{1/2} \leftarrow {}^{2}P_{3/2}$	H,e,H ₂	92 K
[Si II] λ 34.8 μ m ${}^{2}P_{1/2} \leftarrow {}^{2}P_{3/2}$	е	92 K
$[O I] \lambda 63.2 \mu \text{m}^{-3} \text{P}_1 \leftarrow ^3 \text{P}_2$	H,e	228 K
$[O I] \lambda 145.5 \mu \text{m}^{-3} \text{P}_0 \leftarrow^3 \text{P}_2$	H,e	326 K



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Molecular excitation - CO & H₂

ISO observations (Van Dishoeck, 2004, ARA&A, 42, 119) suggest that the contribution from H_2 cooling is more important than estimates based on *IRAS* data: H_2 is among the brightest emission lines in molecular clouds exposed to UV radiation, as for instance in Orion KL (note prominent dust continuuum):





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

fluorescent H₂: not a gas coolant

The rotational excitation of the ground vibrational state of H_2 is usually collisional (i.e. thermal), while higher vibrational states are excited by fluorescence (hence not a cooling process).

 \Rightarrow near-IR rovib H₂ is usually fluorescent, so not a coolant (except some of the lower-level rovib lines, in shocked regions).



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

Dust continuum: not a gas coolant

The most luminous spectral component in the molecular ISM is the mid-IR to far-IR continuum, due to dust. Dust is responsible for transporting the UV radiation of the exciting stars, but is not always coupled to the molecular gas. Collisions are efficient dust heating sources in the denser media only, such as dense SNRs or dense circumstellar disks.





ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

CO rotational-line cooling

The main cooling agents for cold molecular clouds are low-level rotational transitions of CO in the sub-mm. (Figure 2.11 from Tielens 2005).







Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Outline

1 Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density Ionic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Summary: diatomic molecules

In the Born-Oppenheimer approximation the Hamiltonian for a diatomic molecule AB separates, $H = H_{AB} + H_{el}$, where H_{AB} represents the kinetic energy of the nuclei, and H_{el} all the rest. It is also customary to factorize $\phi = \phi_{el}\phi_{AB}$ and obtain ²:

$$H_{\rm el}\phi_{\rm el}=E_{\rm el}(R) \phi_{\rm el},$$

$$H_{\rm AB}\phi_{\rm AB}+E_{\rm el}(R)\ \phi_{\rm AB}=E\ \phi_{\rm AB},$$

and using the centre of mass coordinates, $\phi_{AB} = \phi_{trans}(\vec{X}_{CM})\phi_{intern}(\vec{R}),$

$$-\frac{\hbar^2}{2M}\nabla_{\rm CM}^2\phi_{\rm trans}=E_{\rm trans}\ \phi_{\rm trans}$$

$$-\frac{\hbar^2}{2\mu}\nabla_{\vec{R}}^2\phi_{\rm intern} + E_{\rm el}(R) \ \phi_{\rm intern} = E_{\rm intern} \ \phi_{\rm intern}$$

with $E = E_{\text{intern}} + E_{\text{trans}}$.

²provided additional approximations according to Messiah, Vol. 2, XVIII.64 (p791), but exactly according to Shu I, 28. Probably valid only for processes slower than a rotation period.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

The Laplacian in spherical coordinates,

$$\nabla_{\vec{R}}^{2} = \frac{1}{R^{2}} \left[\frac{\partial}{\partial R} \left(R^{2} \frac{\partial}{\partial R} \right) - \frac{L^{2}}{\hbar^{2}} \right],$$

inspires the decomposition of the wave function for the internal degrees of freedom in spherical armonics,

$$\phi_{\text{intern}}(\vec{R}) = \frac{1}{R} Z_{\text{vib}}(R) Y_{Jm}(\theta, \varphi)$$

,

The 1-D equation for $Z_{vib}(R)$ is a function of $E_{el}(R)$, the energy eigenvalue of the electronic wave function:

$$-\frac{\hbar^2}{2\mu}\frac{d^2 Z_{\text{vib}}}{dR^2} + E_{\text{el}}(R)Z_{\text{vib}} = \left[E_{\text{intern}} - \frac{J(J+1)\hbar^2}{2\mu R^2}\right]Z_{\text{vib}}.$$



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

In the armonic approximation we expand the 1-D equation for $Z_{\rm vib}$ about R_{\circ} , the AB separation in the fundamental state, highlighting the natural frequency $\mu\omega_{\circ} \equiv E''_{el}(R_{\circ})$:

$$-\frac{\hbar^2}{2\mu}\frac{d^2 Z_{\text{vib}}}{dx^2} + \frac{\mu}{2}\omega_{\circ}^2 x^2 Z_{\text{vib}} = E_{\text{vib}}Z_{\text{vib}}$$
$$x \equiv R - R_{\circ}, \text{ and}$$
$$E_{\text{intern}} = E_{\text{el}}(R_{\circ}) + E_{\text{vib}} + E_{\text{rot}},$$

•

 $E_{\rm rot} = J(J+1) \frac{\hbar}{2\mu R_{\circ}^2} \equiv J(J+1)B, \quad B: \text{ rotational constant},$

$$E_{\mathrm{vib}} = (v + \frac{1}{2})\hbar\omega_{\circ}.$$

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

Ionic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation Heating of the ISM

Antonial Antonialism

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

³the only component of \vec{L} that is conserved in diatomic molecules

The quantum numbers that determine the nuclear state of the

independent of J, which is the fundamental hypothesis implicit

in the decomposition $\phi = \phi_{el} \phi_{AB}$. The quantum numbers that

characterise the electronic state of the molecule are Λ , and S.

where Λ corresponds to the projection of L_{tot} along the

internuclear axis³, and S is the total spin.

molecule are J, m, v. Note the rigid body assumption: ϕ_{el} is

Outline

1 Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density Ionic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation



ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Interaction with electromagnetic radiation

The coupling term between charged particles and the electromagnetic field, $\vec{p_i} \cdot \vec{A}(\vec{k} \cdot \vec{x} - wt)^4$, can be expressed through an expansion in $\vec{k} \cdot \vec{x}$ as $H_{\text{int}} = H_{\text{d}} + H_{\text{M}} + H_{\text{Q}}$ (see Shu I,24), for which

$$H_{\rm d} = -\vec{E} \cdot \vec{d}$$
 (zeroth order)

where, for a molecule, $\vec{d} = \vec{d}_{\rm el} + \vec{d}_{\rm nuc}$.

$$H_{\rm M} = -\vec{B}\cdot\vec{M}$$
 (first order)

where the magnetic dipole moment $\vec{M} \propto \vec{L}$, and

$$H_{\rm Q} = -\frac{e}{6}\vec{\nabla}\vec{E}: (3\vec{x}\vec{x} - |\vec{x}|^2\mathbb{I}) \text{ (also order one).}$$

In general $H_{\rm M} > H_{\rm Q}$.

⁴when substituting $\vec{p} \to \vec{p} - \frac{q}{c}\vec{A}$, and neglecting terms in A^2 (OK for the ISM) see Shu I, 21



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Bound-bound transition probabilities and cross-sections

In time-dependent perturbation theory, the rate of radiative excitations $i \rightarrow f$ is:

$$rac{d m{P}_{if}}{dt} \propto \mathfrak{N}(\omega_{if}) \left| \langle \phi_f | m{H}_{ ext{int}}(\omega) | \phi_i
angle
ight|^2$$

In a cubic box where the ocupation number of state \vec{n} is \mathfrak{N} , the density of states is $d^3\vec{n} = Vd^3\omega/(2\pi c)^3$, with a volume $V \to \infty$. The absorption cross-section σ_{if} derives from $P_{if} = N_f/N_i \propto \#$ of absorbed photons ⁵:

$$P_{if} = \int d^3 \vec{n} \, \frac{\mathfrak{N}(\vec{n})}{V} c \, t \, \sigma_{if} \ \Rightarrow \ \frac{dP_{if}}{dt} = \int_0^\infty \sigma_{if} \, c \, \mathfrak{N}(\omega) \, \frac{4\pi\omega^2}{(2\pi)^3 c^2} d\omega$$

Identifying (see Shu I, 22, 23), we obtain

$$\sigma_{if} \propto \left| \langle \phi_f | \mathcal{H}_{int}(\mathbf{w}) | \phi_i \rangle \right|^2 \delta(\omega - \omega_{if}).$$

⁵note optically thin case: $dN_f = -\Gamma N_f dt + \frac{dP_{if}}{dt} N_i dt$.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Oscilator strength

For the electric dipole Hamiltonian, one gets

$$\sigma_{i\!f} = rac{4\pi^2}{3\hbar c} \left| \langle \phi_f | ec{d} | \phi_i
angle
ight|^2 \delta(\omega - \omega_{i\!f}),$$

which is usually expressed in terms of the oscilator strength f_{if} ,

$$\sigma_{if} = \frac{\pi e^2}{m_e c} f_{if} \delta(\nu - \nu_{if}), \text{ con } f_{if} \equiv \frac{4\pi m_e}{3e^2 \hbar} \nu_{if} \left| \langle \phi_f | \vec{d} | \phi_i \rangle \right|^2.$$

For a single electron with position \vec{x} ,

$$f_{if} = rac{2m_e(\omega_{if}\langle f|ec{x}|i
angle)^2}{3\hbar w_{if}},$$

which is roughly the ratio between the vibrational potential energy of the electron and that of the radiated photon.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Relationship with Einstein coeficients

The equation of detailed balance,

$$n_i B_{if} J_{\nu_{if}} = n_f A_{fi} + n_f B_{fi} J_{\nu_{fi}},$$

and the LTE relationships,

$$\frac{n_f}{n_i} = \frac{g_f}{g_i} \exp(-\frac{h\nu_{if}}{kT}), \text{ and } J_{\nu} = B_{\nu}(T), \text{ lead to}$$

$$A_{fi} = rac{g_i}{g_f} (2h
u^3/c^2) B_{if}, \qquad B_{fi} = rac{c^2}{2h
u^3} A_{fi} = (g_i/g_f) B_{if},$$

where the rate of stimulated excitations is related to the oscilator strength⁶:

$$B_{if} = \frac{4\pi^2}{h\nu} \frac{e^2}{m_e c} f_{if}.$$

⁶ using the requirement
$$n_i B_{if} J_{\nu_{if}} = \int d\nu \int d\Omega \sigma_{if} n_i \frac{J_{\nu}}{h\nu} = \int d\nu 4\pi \sigma_{if} n_i \frac{J_{\nu}}{h\nu}$$



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Selection rules

- Electric dipole
 - atoms: $\Delta I = 1$, $\Delta m = 0$.
 - molecules:
 - vibrational-rotational transitions, or rovibrational, ΔJ = ±1, Δm = 0,Δv = ±1, allowed when Λ ≠ 0, ΔJ = 0⁷.
 - electronic transitions, $\Delta \Lambda = 0, \pm 1, \Delta S = 0$
 - electronic-vibrational-rotational transitions (i.e. vibronic transitions): ΔJ = 0, ±1, Δm = 0, ±1 and ΔJ ≠ 0 si Λ = ΔΛ = 0 and if J = 0.

$$\Delta J = \left\{ egin{array}{ccc} +1 &
ightarrow {
m R} {
m branch} \ 0 &
ightarrow {
m Q} {
m branch} \ -1 &
ightarrow {
m P} {
m branch} \end{array}
ight.$$

- magnetic dipole, atoms: $\Delta l = 0$, $\Delta m = 0, \pm 1$.
- electric quadrupole, atoms: $\Delta I = 0, \pm 2, \Delta m = 0, \pm 1, \pm 2,$ rotational transitions in molecules $\Delta J = 0, \pm 1, \pm 2.$



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

 $^{^7}Lambda$ doubling, two states $\pm\Lambda$ for each J. Example: hyperfine structure of the OH Λ doublet at ${\sim}1.7$ GHz.

Selection rules – CO⁸





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

⁸Subaru - IRCS + echelle & X-disperser, Goto et al. 2003, ApJ, 598, 1038
Selection rules – H₂

In the case of H_2 we have $\vec{d} = 0$. Moreover the fundamental state is $\vec{L} = 0$, and $\vec{S} = 0$, so that $\vec{M} = 0$ and all low-energy transitions for H_2 are quadrupolar.

Note that the antisymmetry of the nuclear wave function implies that the state J = 1 (J odd) is triplet (**Ortho** H₂, $I = S_{nuclear} = 1$), while J = 0 (J even) is singlet (**Para** H₂, I = 0). In H₂ the exclusion principle ⁹ forbids $\Delta J = 1$, unless the transition involves a change in spin state. The spin transitions can only occur through the exchange of protons in collisions. Radiative transitions between spin states can occur, but at a rate corresponding to the quadrupolar transitions in the Hamiltonian of the deviations to the Born-Oppenheimer approximations.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

⁹the requirement that the wave function be antisymmetric



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

In the ISM, **Ortho** and **Para** H_2 are effectively different molecules. The distinction extends to all molecules that contain H_2 radicals.

Rovibrational transitions between an upper level ¹ and a lower level ² are written $(v_1 - v_2)O(J_2)$, when $J_2 - J_1 = -2$, $(v_1 - v_2)Q(J_2)$ when $J_2 - J_1 = 0$, $(v_1 - v_2)S(J_2)$ when $J_2 - J_1 = +2$.

H₂ energy levels



Figure 1 Total energy as a function of internuclear separation for low-lying states of H₂. Horizontal dashed lines denote total energy of two H atoms at infinite separation (both atoms in ground state; one atom in ground state and the other in the first excited states).

 $\begin{array}{l} \Lambda = 0, 1, 2, ... \\ \Leftrightarrow \Sigma, \Pi, \Delta \\ \text{levels in } E_{\text{el}}(R) \\ X, B, C, D. \end{array}$

The energy separation between levels *B* and *C* with respect to *X* are \sim 11 eV and 14 eV.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Dissociation of H₂

The collisional dissociation of H₂,

$$\begin{array}{ccc} He & He \\ H_2 + & H \\ e & 2H + & H \\ e & e \end{array},$$

requires temperatures of order the binding energy of H₂, 4.48 eV/k = 52000 K¹⁰. But an H₂ cloud is more likely to be photo-dissociated before it reaches 50 000 K. The dissociation continuum of H₂ (i.e. free states of H+H) starts 14.7 eV above the fundamental level of H₂ (i.e. X¹ Σ , para-H₂), while the ionization continuum (the ejection of one electron) starts at 15.4 eV. Both continua, ionising and dissociating, lie above 13.6 eV. In the context of a molecular cloud surrounded by H I, the H₂ dissociating radiation is completely absorbed by the Lyman continuum of H I.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

¹⁰the law of mass-action does not change this result because the number of continuum states accessible to both products and reactants is similar, in contrast with ionisation, which is described by the Saha equation

Dissociation of H₂

The likeliest mechanism for H_2 dissociation is absorption to the robrivational bands of excited electronic states (e.g. B o C), and

- subsequent de-excitation to dissociated levels of lower energy (e.g. b³Σ), or
- 2 decay to the vibrational continuum of electronic state X (with v ≥ 14), or
- decay to excited levels of the fundamental state and subsequent absorption to the dissociating continuum (one order of magnitude slower than mechanisms 1) and 2) , Stecher& Williams 1967, ApJ 149, L29)



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

IR emission from H₂

Since the rate of radiative transitions within the ground state electronic level of H_2 is quadrupolar, the distribution of ro-vibrational population is dominated by collisions between H_2 molecules.

If $kT \gg \Delta E_{\rm FIR}$, the typical energy difference between rotational levels of ground state H₂, the occupation numbers of the robrivational levels in the ground state electronic state follow LTE:

 $N(v, J) \propto \omega(J) \exp(-B(J(J+1))/kT),$

in which $\omega(J) = 3 \times (2J + 1)$ for ortho- H_2 , and $\omega(J) = 1 \times (2J + 1)$ for para- H_2 . This situation (i.e. hot molecular gas) is typical of shocks in the ISM. In the presence of an intense UV field the IR emission of H_2 is fluorescent (Black & Van Dishoeck, 1987, ApJ, 322, 412). Fluorescence dominates if $kT \ll \Delta E_{\rm FIR}$. The absorption of one UV photon excites H_2 to an electronic state, leading to photodissociation (with a 15% probability), and to an IR cascade in all other cases.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

IR emission from H₂

Example excitation diagram for the Helix nebula.



IR emission from H_2 , ζ Oph







ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Outline

1 Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density Ionic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation



ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Heating of the neutral ISM

 $\rm H_2$ is an important source of UV opacity, as shown in the FUSE spectra of HD210839 and HD154368 (Rachford et al., 2002, ApJ, 577, 221).





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Heating mechanisms

1 Photoionisation:

$$A + h\nu \rightarrow A^+ + e_{\mu}$$

in the case of the neutral ISM, photoionisation heating is dominated by C^+ , with an ionization potential of 11.3 eV. The average kinetic energy of the ejected electron is

$$\frac{3}{2}kT_{i} = \frac{\int_{\nu_{i}}^{\infty} (J(\nu)/h\nu)\alpha_{i}(\nu)h(\nu-\nu_{i})d\nu}{\int_{\nu_{i}}^{\infty} (J(\nu)/h\nu)\alpha_{i}(\nu)d\nu}$$

2 Photodissociation:

 $AB + h\nu \rightarrow A + B + KE$,

dominated by H_2 photodissociation.

3 collisional de-excitation (in particular for H₂, when n(H₂) > 10⁴ cm⁻³).



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

4 photoelectric effect on dust grains.

The grain heating efficiency, ϵ_{grain} , defined as the ratio between output gas heating and input UV radiation (Tielens Chap. 3.3.1), is

$$\epsilon_{\text{grain}} = \mathbf{Y} imes rac{h
u - W - \phi}{h
u}$$

•

where the electron escape probability *Y* is the ratio between the e^- - e^- mean free path $I_e \approx 10 R$, and the absorption depth of UV radiation, $I_a \approx 100 R$: $Y \approx I_e/I_a$. *W* is the work function of the grain and $\phi = Z \times e^2/a$ is the surface potential for a grain of size *a*. The ionization potential of a grain is thus $I = W + \phi$.

The heating efficiency is larger for smaller grains, since $l_a \approx l_e$. The ejection of photoelectrons from grains is the dominant heating source in the neutral ISM exposed to FUV radiation, as in the case of PDRs. Photoionisation is dominant in the case of the ionised gas.

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis Ion-molecule chemistry Observational validation Chemical fractionation

The total photoelectric heating rate from grains per unit volume, nΓ_{grain}, can be written:

$$n\Gamma_{\text{grain}} = \int_{a_{\min}}^{a_{\max}} n(a) da \sum_{i} \int_{\nu_{i}}^{\nu_{H}} \frac{J_{\nu}}{h\nu} \alpha_{\text{grain},i}(\nu) E_{\text{kin}}(a,i) d\nu,$$

where *a* is grain diameter, and *i* runs over grain ionisation stage, and $h\nu_{\rm H} =$ 13.6 eV.

 in addition to photo-electric heating, grains can also heat the gas through inelastic collisions if the solid state temperature is larger than the gas temperature.

b heating from cosmic rays:

Important for molecular phase. Cosmic rays, as well as X-rays, ionise H_2 (IP=15.4 eV):

$$H_2 + (p, e^-, \alpha, X) \rightarrow H_2^+ + e^- + (p, e^-, \alpha, X)',$$

followed by the formation of H_3^+ :

$$\mathrm{H_2^+} + \mathrm{H_2} \rightarrow \mathrm{H_3^+} + \mathrm{H}.$$

An important mechanism for the destruction of H_3^+ is its dissociative recombination,

$$\mathrm{H}_3^+ + \textbf{e}
ightarrow \left\{ egin{array}{c} \mathrm{3H} \\ \mathrm{H}_2 + \mathrm{H} \end{array}
ight.$$

Since both the formation and the destruction of H_3^+ are exothermic (the energy released into the ISM in the process is 11 eV), the H_3^+ cycle is a heating source of the ISM.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

6 Turbulent heating.

Important in molecular clouds at \sim 10 K, without UV field. In a turbulent flow the kinetic energy density will trickle down a hierarchy of scales, until it is dissipated by diffusion at the smallest scale fixed by viscosity. The timescale for downwards energy transfer at scale / can be approximated with the crossing time at the velocity disperson v. Thus the rate of heating per unit volume is (Tielens 3.8),

$$n\Gamma_{\rm turbulence} = \frac{1}{2} \frac{n \ m_{\rm H} v^2}{l/v}$$

The observations of the energy density and velocity dispersion at a given scale *I* leads to an estimate of turbulent heating.

The energy densities as a function of scale can be calculated assuming a constant energy transfer rate: $v^2/(l/v)$ is constant.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation







Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation





ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Steady state temp. for HI





ISM

Two phases for HI





ISM

Molecule formation

Outline

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical densit Ionic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis

Ion-molecule chemistry Observational validation Chemical fractionation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry Observational validation Chemical fractionation

Molecule formation

The encounter of two atoms *A* and *B*, initially unbound, will lead to molecular synthesis if:

- 1 bound states of *AB* exist.
- 2 There is a mechanism for molecular stabilization faster than the lifetime of the AB^{*} state ¹¹, of order ~ 10⁻¹³ s (≈ 1 vibrational period).

Radiative dissipation of the energy excess of AB^* is improbable: even for electric dipole decay rates, about 10^8 s^{-1} in the UV, only 1 out of 10^5 collisions will produce a molecule. The abundance of H₂ might raise the low probability of condensation via collisions, but the electric dipole transitions of H₂ are fobbiden. H₂ cannot associate radiatively. The H₂ half-life in the diffuse interstellar UV radiation field is of ~ 300 yr \Rightarrow some other mechanism must exist for efficient H₂ formation.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry Observational validation Chemical fractionation

¹¹the molecule is initially bound at a level with energy equal to that of the colliding pair

Grain catalysis of the formation process of H₂

The adsorption of hydrogen atoms onto the surface of dust grains allows raising the available timescale for the A + B encounter, and the grain acts as stabilizing agent by absorbing the excess energy of AB^* .

Once H_2 is formed, its 4.48 eV of binding energy (the first bound-free transition is at 2500 Å) eject H_2 from the grain surface, and the excess of energy over the grain work function is distributed over the molecular degrees of freedom, leaving the H_2 molecule in an excited state.

IR transitions from the formation pumping of excited levels of H₂ were observed by Burton et al. (2002, MNRAS, 333, 721) in M 17: the line (6-4)O(3) at 1.73 μ m cannot be collisionally excited since its superior level is at $\Delta E/k \sim 30\,000$ K above ground, yet the intensity of (6-4)O(3) relative to other H₂ lines does not match fluorescent models.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry Observational validation Chemical fractionation

Outline

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical densit Ionic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis Ion-molecule chemistry Observational validation

Chemical fractionation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry

Observational validation Chemical fractionation

Molecule formation - the role of H₃⁺

• H_3^+ controls the ionisation stage of the molecular ISM. The H_3^+ cycle with corresponding rates is:

$$\begin{aligned} \mathrm{H}_{2} + (\boldsymbol{p}, \boldsymbol{e}^{-}, \alpha, \boldsymbol{X}) &\stackrel{\zeta}{\rightarrow} \mathrm{H}_{2}^{+} + \boldsymbol{e}^{-} + (\boldsymbol{p}, \boldsymbol{e}^{-}, \alpha, \boldsymbol{X})', \\ \mathrm{H}_{2}^{+} + \mathrm{H}_{2} \stackrel{\eta_{1}}{\rightarrow} \mathrm{H}_{3}^{+} + \mathrm{H}, \\ \mathrm{H}_{3}^{+} + \boldsymbol{e} \stackrel{\eta_{2}}{\rightarrow} \mathrm{H}_{2} + \mathrm{H}, \end{aligned}$$

where ζ is the ionization rate of H₂, in s⁻¹. We see that

$$\zeta n_{
m H_2} = \eta_1 n_{
m H_2} n_{
m H_2^+}, \; {
m and} \;$$

$$\eta_2 n_{\mathrm{H}_3^+} n_e = 2\zeta n_{\mathrm{H}_2}$$

With canonical values for the rates η_i , one gets $n_{{
m H}_3^+} \gg n_{{
m H}_2^+}$, and the requisite for neutrality $n_{
m e} \approx n_{{
m H}_3^+}$ allows concluding $n_{
m e}/n_{{
m H}_2} \approx 10^{-8}$.



ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry

Observational validation Chemical fractionation

Ion-molecule chemistry

• In ion-molecule chemistry, the standard model for interstellar chemistry, H_3^+ is cornerstone to all reaction cycles: The cycle of an element X (anything short of He, O₂,...) is initiated by H_3^+ through

 $H_3^+ + X \rightarrow XH^+ + H_2$, and subsequently,

 $XH^+ + Y \rightarrow XY^+ + H.$

Example: production of H₂O and OH,

 $\rm H_3^+ + O \rightarrow OH^+ + H_2, \mbox{ followed by the sequence,}$ $\rm OH^+ + H_2 \rightarrow OH_2^+ + H,$ $\rm OH_2^+ + H_2 \rightarrow OH_3^+ + H,$

and finally one dissociative recombination

$$\mathrm{OH_3}^+ + \mathrm{e} \rightarrow \left\{ \begin{array}{c} \mathrm{H_2O} + \mathrm{H} \\ \mathrm{OH} + \mathrm{H_2} \end{array} \right.$$



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry

Observational validation Chemical fractionation

The full story

• Molecular formation: realistic cycles (Sternberg & Dalgarno 1995, ApJSS, 99, 565).





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry

Observational validation Chemical fractionation

ISM

Outline

1 Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical densit Ionic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry

Observational validation

Molecular formation - validation of the ion-molecule chemical model

 H_3^+ , the pilar of chemical cycles in the ISM, was detected by Geballe & Oka (1996, *Nature*, 384, 334; Geballe 2000, *Phil.Trans.R.Soc.Lond.A*, 358, 2503), with an abundance close to that expected from the ionization fraction.

 H_3^+ does not have electric dipole moments, and is even worse than H_2 : it is deprived of excited electronic states (i.e. no UV transitions). It only has two vibration modes. One is symmetrical, ν_1 , that does not induce electric dipole moment and does not radiate. Only the asymmetrical mode ν_2 induces an electric dipole.





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

lon-molecule chemistry Observational validation

Chemical fractionation

ISM

The production of H₃⁺

$$\begin{split} \mathrm{H}_{2} + (\pmb{p}, \pmb{e}^{-}, \alpha, \pmb{X}) &\xrightarrow{\zeta} \mathrm{H}_{2}^{+} + \pmb{e}^{-} + (\pmb{p}, \pmb{e}^{-}, \alpha, \pmb{X})', \\ \mathrm{H}_{2}^{+} + \mathrm{H}_{2} &\xrightarrow{\eta_{2}} \mathrm{H}_{3}^{+} + \mathrm{H}, \end{split}$$

is balanced by its destruction, mainly through

$$H_3^+ + CO \xrightarrow{\eta_{CO}} COH^+ + H_2$$

equating rates, one gets

$$2\zeta n_{\rm H_2} = \eta_{\rm CO} n_{\rm H_3^+} n_{\rm CO}$$

in which $\eta_{\rm CO} = 1.8 \ 10^{-9} \ {\rm s}^{-1}$, and $\zeta = n_{\rm p,X,\alpha,e^-} \ \eta_1 \approx 3 \ 10^{-17} \ {\rm s}^{-1}$, gives the product of the cosmic ray density times the ionisation rate of H₂. In general the relative abundances of CO and H₂ are constant in the diffuse ISM (model: Lee et al., 1996, A&AS, 119, 111, observations: Federman et al. 1980, ApJ 242, 545), $n_{\rm CO}/n_{\rm H_2} = 1.5 \ 10^{-4}$, so that

 $n_{\mathrm{H_3^+}} pprox 10^{-4} \ \mathrm{cm^{-3}}, \ \ \mathrm{independent} \ \mathrm{of} \ n_{\mathrm{H_2}}$

For typical densities $n_{\rm H_2} \approx 10^4 {\rm cm}^{-3}$, we have $N_{\rm H_3^+} \approx 10^{-9} N_{\rm H_2}$!!



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Detection of H⁺₃, **Oka & Geballe**





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis Ion-molecule chemistry Observational validation





ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis Ion-molecule chemistry

Observational validation





Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis

Ion-molecule chemistry

Observational validation

Outline

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical densit Ionic & atomic cooling of the ISM Molecular excitation & molecular structure Atomic/molecular coupling with radiation Heating of the ISM

2 Molecule formation

Grain catalysis Ion-molecule chemistry Observational validation Chemical fractionation



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Chemical fractionation: CO

The reaction

$$^{13}\mathrm{C}^+ + {}^{12}\mathrm{CO} \rightleftharpoons_{k_2}^{k_1} {}^{12}\mathrm{C}^+ + {}^{13}\mathrm{CO} + \Delta E,$$

is slightly exothermic: $\Delta E/k = 35$ K. Consequently, at low temperatures ($\lesssim 35$ K), the isotopic ratio ${}^{12}C/{}^{13}C$ is different in CO than in the rest of the ISM - there is fractionation in the C isotopes.



ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Chemical fractionation: theory

At constant *P* and *T*, minimizing the free energy *G* gives $\sum_i \nu_i \mu_i = 0$, and allows relating the concentration of each species involved in the reaction. For reactions in the gas phase, we can use the chemical potential of ideal gases, $\mu_i = kT \ln [n_i/\mathcal{Z}_i(T)]$, with the partition function

$$\mathcal{Z}_i(T) = \frac{h}{\sqrt{2\pi m_i \, k \, T}} \sum_a g_i^a e^{-\xi_a/kT}$$

where the energies associated with the internal degrees of freedom are referred to the ground state. We get the law of mass-action (e.g. Shu I, 7),

$$\prod_i n_i^{\nu_i} = \prod_i \mathcal{Z}_i^{\nu_i}.$$



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation

ISM
Application to CO

For CO fractionation,

$$\frac{n({}^{13}\mathrm{C}^{+}) n({}^{12}\mathrm{CO})}{n({}^{12}\mathrm{C}^{+}) n({}^{13}\mathrm{CO})} = e^{-\frac{1}{kT} \left(\xi({}^{12}\mathrm{CO}) - \xi({}^{13}\mathrm{CO})\right)} = \exp\left(-\frac{\Delta E}{kT}\right)$$

where the ground state energy is ξ (¹²CO).

For CO, the rates of reactions in both directions are fast, $k_1 = k_2 = 2 \ 10^{-10} \ \mathrm{cm}^3 \ \mathrm{s}^{-1}$, and for cold clouds with $n(H_2) \approx 10^4 \ \mathrm{cm}^3$ and $n_{\mathrm{CO}}/n_{\mathrm{H_2}} \approx 10^{-4}$, the relaxation time is $\tau_{\mathrm{frac}} \lesssim 1000 \ \mathrm{yr}$. Thus CO is enriched in ¹³C, at the expense of the rest of the molecules.

The work function of a dust grain for the adsorption of CO is $\Delta H \sim 900$ K, and for $T \leq 900$ K the probability of adsorption for grain-CO collisions is 1. Besides the vapor pressure of CO is very high compared to typical ISM pressures (Leger, 1983, A&A, 123, 271). CO is expected to be substantially (of order 50%) condensed onto icy CO mantles on the surface of grains (Whittet et al. 1985, MNRAS, 216, 45).



ISM

Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation

Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation



We see that the isotopic ratio ${}^{12}C/{}^{13}C$ will be different in the dust than in the gas phase. Additionally the C ratio in any given molecule of the gas phase will depend on the details of the synthesis reactions. The position taken

by CO in relation to C⁺ will determine the C ratio in the product molecule. This result bears on the study of the chemical evolution of galaxies, as well as on studies of molecular cloud structure based on the optically thin isotopes. Chemical fractionation competes with the process of selective photodissociation. ¹³CO is dissociated by interstellar UV radiation, penetrating to deeper lying layers in the clouds than for the main isotope, ¹²CO.



Thermal Equilibrium of the ISM

Cooling of the ISM: detailed balance, critical density

lonic & atomic cooling of the ISM

Molecular excitation & molecular structure

Atomic/molecular coupling with radiation Heating of the ISM

Molecule formation

Grain catalysis lon-molecule chemistry Observational validation Chemical fractionation